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RESEARCH ARTICLE

AC BACK SURFACE RECOMBINATION IN $n^+ - p - p^+$ SILICON SOLAR CELL: EFFECT OF TEMPERATURE

Denise Kabou, Mamadou Lamine Ba, Mamour Amadou Ba, Gora Diop, El Hadji Sow, Oulimata Mballo and Gregoire Sissoko

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Abstract

Excess minority carrier's diffusion equation in the base of monofaciale silicon solar cell under frequency modulation of polychromatic illumination is resolved. Using conditions at the base limits involving recombination velocities S_f and S_b , respectively at the junction (n^+/p) and back surface (p^+/p), the expression of the excess minority carriers' density $\delta(T, \omega)$ is determined. The density of photocurrent $J_{ph}(T, \omega)$ is represented according to the recombination velocity at the junction for different temperature values. The expression of the ac back surface recombination velocity S_b of minority carriers is then deduced depending on frequency of modulation, temperature, electronic parameters (D) and the thickness of the base. Bode and Nyquist diagrams are used to analyze, both diffusion coefficient and back surface recombination of excess minority carriers submitted to Umklapp process.

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Introduction:-

The solar cell AC response induced by modulated optical or electrical signal have been used by previous works[1], to determine either the electronics parameters (τ , L , D , S_b)[2],[3], [4]; [5], [6], [7]or the electrical one($Z(\omega)$, $G(\omega)$, $C(\omega)$, R_s , R_{sh})[8], [9], [10]; [11], [12], [13],[14].The external conditions under which the solar cell is placed are taken into account, in particular, the magnetic field[15], [16], the flow of irradiation by charged particles[17], [18], [19], as well as the temperature[20], [21], [22], [23] and the illumination wavelength[24], [25].Intrinsic parameters were considered such as, the doping rate of and the geometric parameters[24], [25], [26], [27], [28], [29], [30], [31].

Our work is interested in the $n^+ - p - p^+$ solar cell, under composite illumination in frequency modulation and placed under different temperatures (T), in order to study the AC back surface recombination.

Thus the diffusion coefficient of minority carriers in excess $D(\omega, T)$ is represented and analyzed.

The expression of the density of photogenerated carriers is given by the solution of the diffusion equation in the base of the solar cell.

The photocurrent density is then calculated as a function of:

1. surface recombination velocities at the junction (S_f and at the rear (S_b)[32], [33], [34][35],
2. the temperature (T)
3. the illumination modulation frequency (ω) .

Corresponding Author:- Grégoire Sissoko

Address:- Laboratory of Semiconductors and Solar Energy, Physics Department, Faculty of Science and Technology, University Cheikh Anta Diop, Dakar, Senegal.

The back recombination velocity expression is deduced as a function of both the temperature (T) and the frequency (ω) of illumination modulation. It is finally analyzed using, the Umlapp process, showing the thermal agitation of charge carriers [23], [36], [37] and the spectroscopy method of Bode and Nyquist diagrams (amplitude and phase). [11], [17], [27], [35], in order to define the electrical circuit model representation. The real and imaginary parts are plotted versus temperature and modeled through mathematical relationships.

Theory:

The structure of the $n^+ - p - p^+$ monofaciale silicon solar cell [38], [39] under polychromatic illumination, in frequency modulation, is given by figure 1.

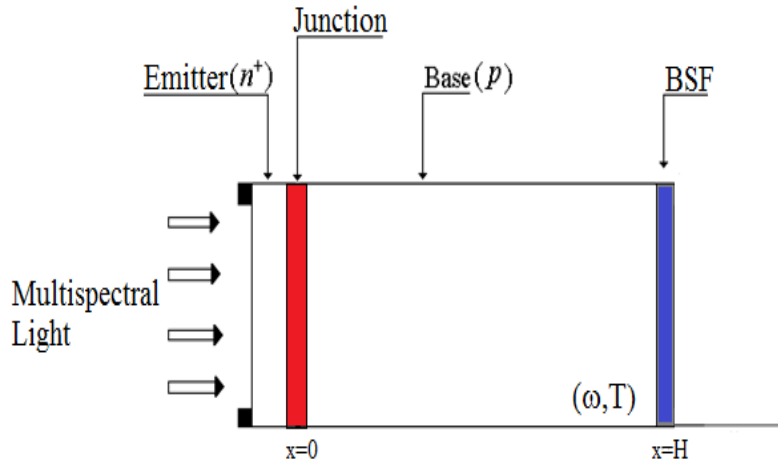


Figure 1:- Structure of monofaciale solar cell.

The excess minority carriers' density $\delta(x, t)$ generated in the base of the solar cell obeying to the continuity equation at T temperature, under polychromatic illumination in frequency modulation, is given by [3], [7], [24]:

$$D(\omega, T) \times \frac{\partial^2 \delta(x, t)}{\partial x^2} - \frac{\delta(x, t)}{\tau} = -G(x, \omega, t) + \frac{\partial \delta(x, t)}{\partial t} \quad (1)$$

The expression of the excess minority carriers' density is written, according to the space coordinates (x) and the time t, as:

$$\delta(x, t) = \delta(x) \cdot e^{-j\omega t} \quad (2)$$

- Carriers generation rate $G(x, t)$ is given by the relationship [39]:

$$G(x, t) = g(x) \cdot e^{-j\omega t} \quad (3)$$

With:

$$g(x) = \sum_{i=1}^3 a_i \cdot e^{-b_i \cdot x} \quad (4)$$

- The depth in the base is represented by x.

- Coefficients a_i et b_i are obtained from tabulated values of radiation in AM 1.5 conditions.

- $D(\omega, T)$ is the complex diffusion coefficient of excess minority carrier in the base at T temperature.

Its expression is given by the relationship [3], [7], [24]:

$$D(\omega, T) = D(T) \times \left(\frac{1 - j \cdot \omega^2 \cdot \tau^2}{1 + (\omega \tau)^2} \right) \quad (5)$$

$D(T)$ is the temperature-dependent diffusion coefficient given by Einstein's relationship [20], [22]:

$$D(T) = \frac{\mu(T) \cdot K_b \cdot T}{q} \quad (6)$$

T is the temperature in Kelvin, K_b is the Boltzmann constant:

$$K_b = 1.38 \times 10^{-23} \text{ m}^2 \cdot \text{Kg} \cdot \text{S}^{-1} \cdot \text{K}^{-1}$$

The mobility coefficient [20] for electrons, expressed according to the temperature, is given by:

$$\mu(T) = 1.43 \cdot 10^{19} T^{-2.42} \quad (7)$$

By replacing equations (2) and (3) in equation (1), the continuity equation for the excess minority carriers' density in the base is reduced to the following relationship:

$$\frac{\partial^2 \delta(x, \omega)}{\partial x^2} - \frac{\delta(x, \omega)}{L^2(\omega, T)} = -\frac{g(x)}{D(\omega, T)} \quad (8)$$

$L(\omega, T)$ is the complex diffusion length of excess minority carriers in the base given by :

$$L(\omega, T) = \sqrt{\frac{D(\omega, T)\tau}{1 + j\omega\tau}} \quad (9)$$

τ is the excess minority carriers lifetime in the base.

1) The solution of equation (8) is:

$$\delta(x, \omega, T) = A \cdot \cosh\left[\frac{x}{L(\omega, T)}\right] + B \cdot \sinh\left[\frac{x}{L(\omega, T)}\right] + \sum K_i \cdot e^{-b_i \cdot x} \quad (10)$$

$$\text{With } K_i = -\frac{a_i \cdot [L(\omega, T)]^2}{D(\omega, T)[L(\omega, T)^2 \cdot b_i^2 - 1]}$$

$$\text{and } (L(\omega, T)^2 \cdot b_i^2 \neq 1) \quad (11)$$

Coefficients A and B are determined through the boundary conditions:

- At the junction ($x = 0$)

$$D(\omega, T) \left. \frac{\partial \delta(x, T)}{\partial x} \right|_{x=0} = S_f \cdot \left. \frac{\delta(x, T)}{D(\omega, T)} \right|_{x=0} \quad (12)$$

- On the back side in the base ($x = H$)

$$D(\omega, T) \left. \frac{\partial \delta(x, T)}{\partial x} \right|_{x=H} = -S_b \cdot \left. \frac{\delta(x, T)}{D(\omega, T)} \right|_{x=H} \quad (13)$$

S_f and S_b are respectively the recombination velocities of the excess minority carriers at the junction and at the back surface. The recombination velocity S_f reflects the charge carrier's velocity crossing the junction, in order to contribute to the photocurrent. It is then related to the external load which imposes the solar cell operating point [30], [32], [33]. It has an intrinsic component which produces the carrier losses associated with the shunt resistor in the solar cell electrical equivalent model [17], [19] [34]. The excess minority carrier recombination velocity S_b on the back surface is associated with the presence of the p^+ layer which produces an electric field for returning back the charge carriers' across the junction [38], [40], [41], [42].

Results and Discussions:-

Diffusion Coefficient: Bode and Nyquist Diagrams at Different Temperatures:

The amplitude and phase of the diffusion coefficient under different temperature are represented versus frequency through figures 2 and 3.

For a given temperature, the effective diffusion coefficient is maximum and constant when the frequency is low (steady state). In a dynamic frequency regime, the repeated excitations of excess minority carriers lead to a relaxation problem that slows their diffusion and brings out a negative phase of the diffusion coefficient. Indeed, the diffusion of minority carriers decreases with an increase in temperature. The effective diffusion coefficient is more sensitive to temperature in a quasi-static regime.

The Nyquist diagram is shown in figure 4.

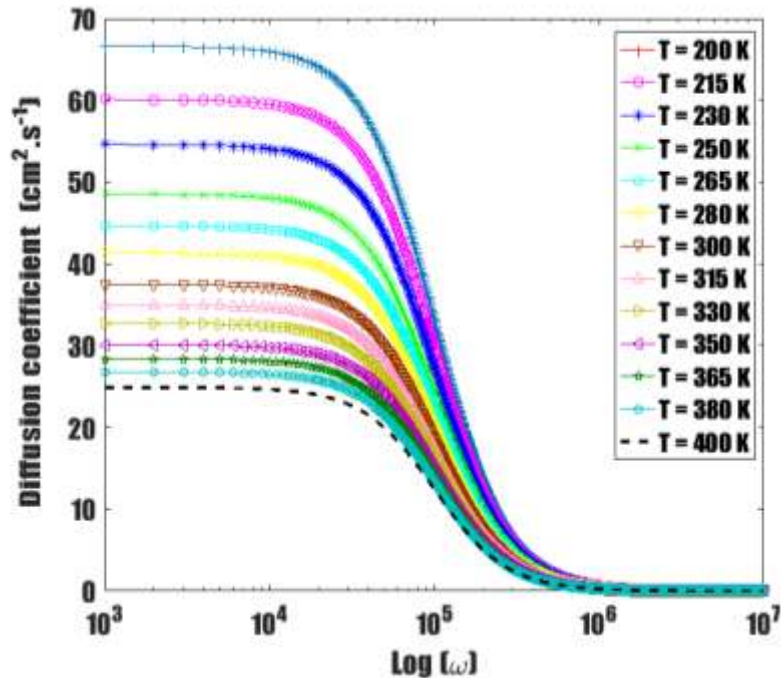


Figure 2:- Diffusion coefficient versus frequency for different temperatures.

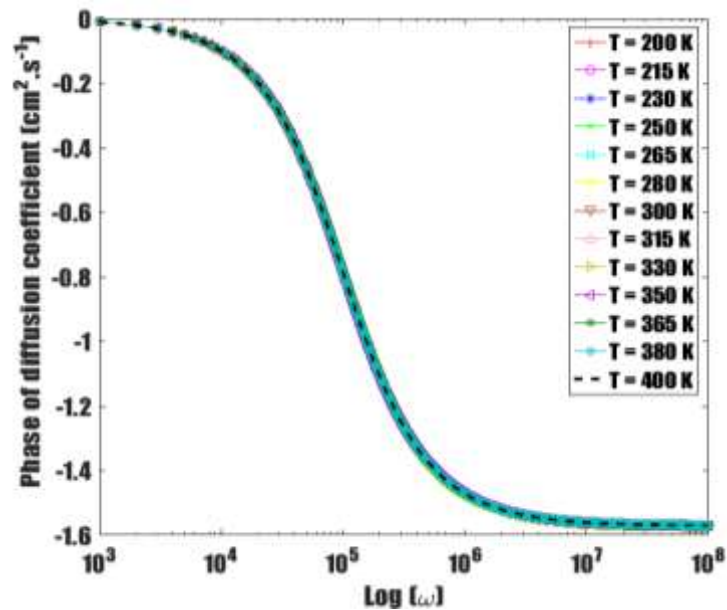


Figure 3:- Diffusion coefficient phase versus frequency for different temperatures.

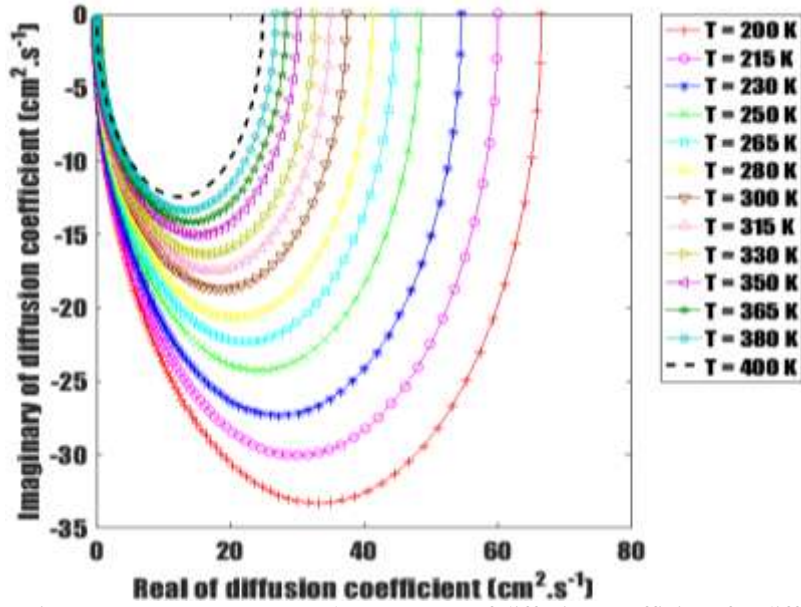


Figure 4:- Imaginary component versus real component of diffusion coefficient for different temperatures.

We find that the radius of the semicircles decreases according to the temperature with a shift from the center of the circles towards the origin of the axes. The semi-circle has no shift at the origin of axes, indicates then a resistor(R_p) in parallel with a capacitor(C)[11], [43], [45], [46] and so gives rise to a single time constant ($R_p.C$), temperature dependent. The investigation of the half-circle radius, thus suggests electrical parameters characteristic leading to equivalent electric model[11],[27], [35], [28].

Minority carrier’s density for different frequencies and temperatures:

The figures. 5, 6 and 7, give the profile of the density of minority carriers in excess with the depth of the base, for different frequencies and at a given temperature.

Figure 5 shows an amplitude of density of minority carriers in excess that is increasing with temperature, and is insensitive to low frequencies (steady state in amplitude of effective diffusion coefficient)

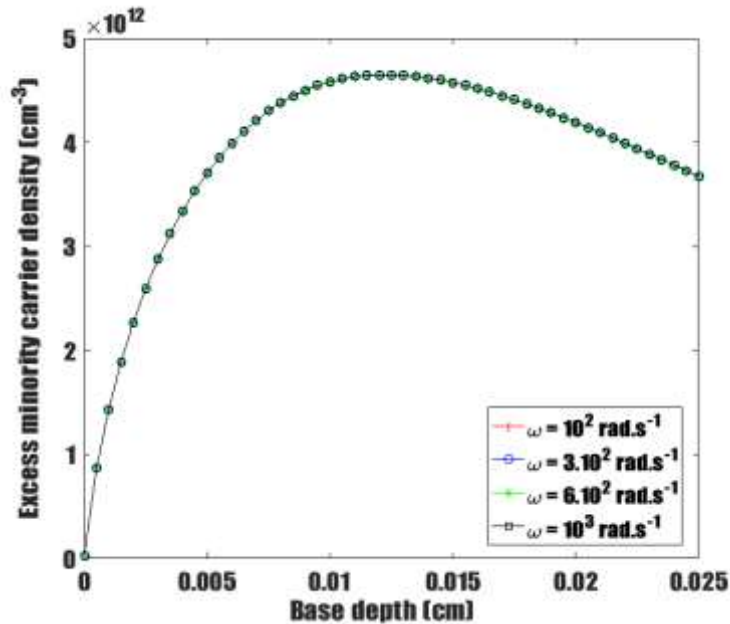


Figure 5:- Excess minority carrier’s density versus base depth for different frequencies (T=200 K)

At low temperatures (Fig. 6), the amplitude of the density of minority carriers in excess decreases due to the low thermal agitation and remains sensitive to variations in frequency (region of abrupt variation of the diffusion coefficient).

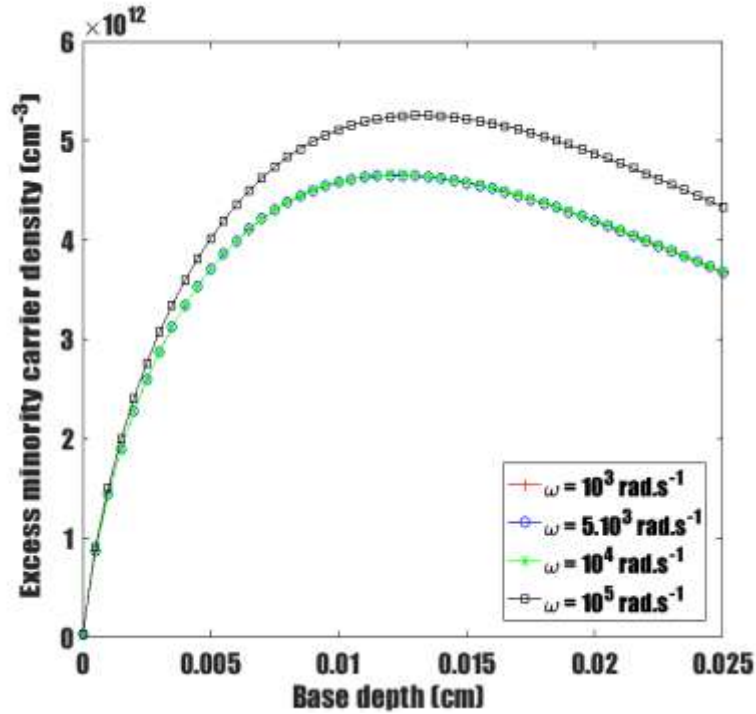


Figure 6:- Excess minority carrier’s density versus base depth for different frequencies (T=200 K).

When temperature increases, the density of excess carriers undergoes thermal agitation, by increasing amplitude (Umklap process) and variation in frequency (in stationary zone) is insensitive (Fig. 7) [36], [37].

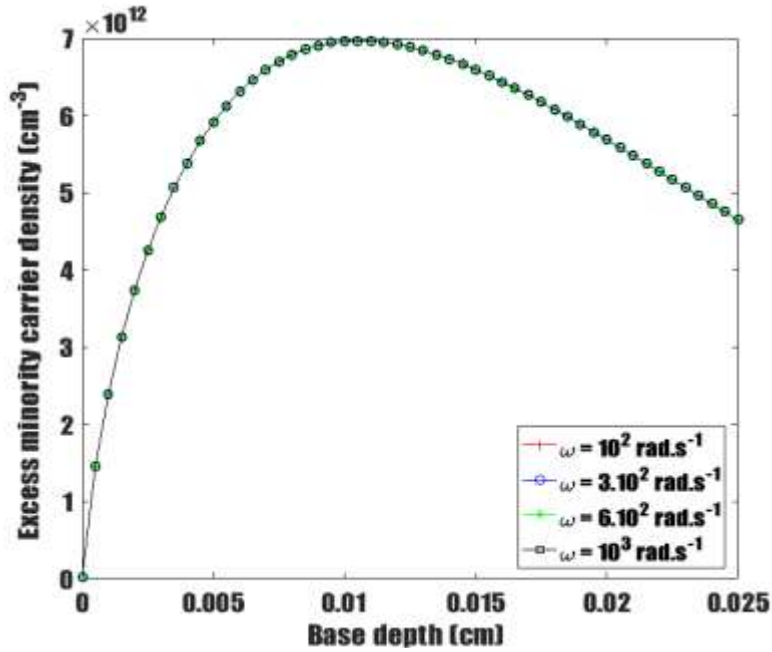


Figure 7:- Excess minority carrier’s density versus base depth for different frequencies (T=300 K).

In Figure 8, in addition to the high temperature requiring greater thermal agitation of the minority carriers, the increase in frequency on the other hand decreases the amplitude of the density of the carriers with the base depth [21], [23], [36], [37].

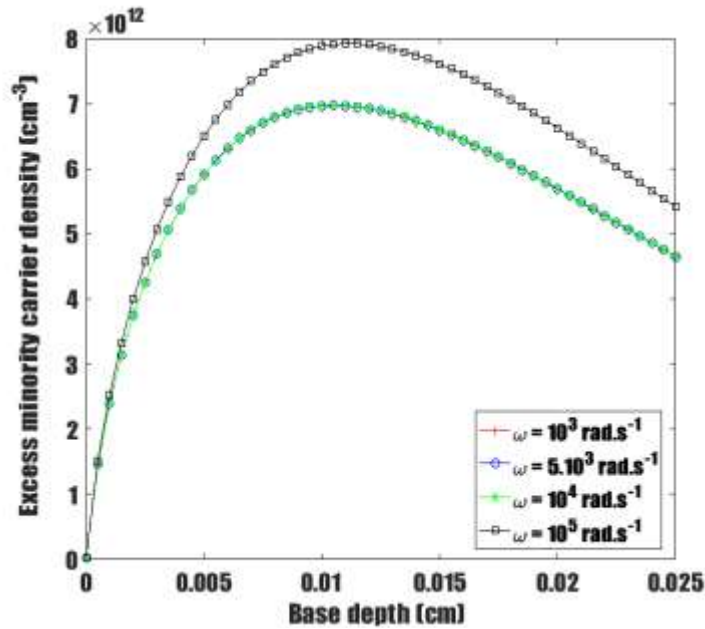


Figure 8:- Excess minority carrier’s density versus base depth for different frequencies (T=300 K)

Photocurrent:

The density of photocurrent at the junction is obtained from the density of minority carriers in the base and is given by the following expression:

$$J_{ph}(Sf, Sb, \omega, T) = qD(\omega, T) \left. \frac{\partial \delta(x, Sf, Sb, \omega, T)}{\partial x} \right|_{x=0} \tag{14}$$

Where q is the elementary electron charge.

Figure 5 shows ac photocurrent versus the junction surface recombination velocity for different temperatures.

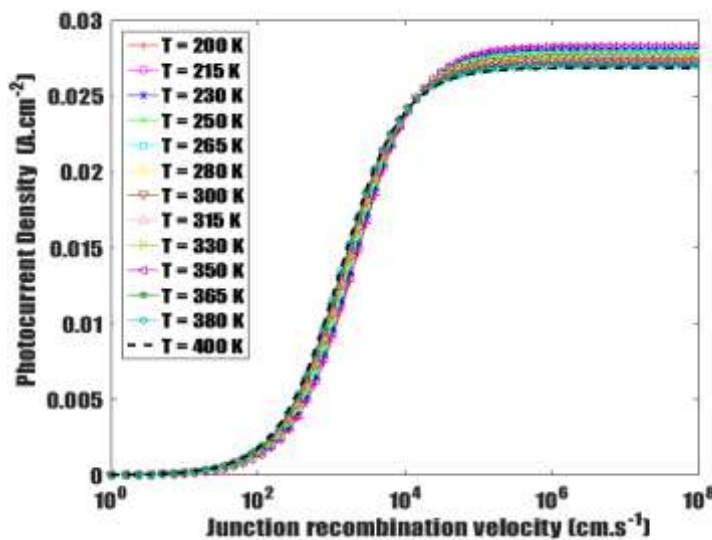


Figure 9:- Photocurrent density versus junction surface recombination velocity for different temperatures ($\omega=10^5$ rad/s).

Deduction of the Sb (ω , T) expression:

The representation of photocurrent density (**figure. 9**) according to the junction recombination velocity of minority carriers shows that, for very large Sf, a bearing sets up and corresponds to the short-circuit current density (Jphsc). So, in this junction recombination velocity interval, it comes [30], [32], [33], [34]:

$$\left. \frac{\partial J_{ph}(Sf, Sb, \omega, T)}{\partial Sf} \right|_{Sf \geq 10^5 \text{ cm.s}^{-1}} = 0 \quad (15)$$

The solution of the above equation leads to expressions of the ac recombination velocity in the back surface, given as:

$$Sb1(\omega, T) = \frac{D(\omega, T)}{L(\omega, T)} \sum_{i=1}^3 \frac{L(\omega, T) b \left(\exp(-bi.H) - \cosh\left(\frac{H}{L}\right) \right) - \sinh\left(\frac{H}{L}\right)}{L(\omega, T) \cdot bi \cdot \sinh\left(\frac{H}{L(\omega, T)}\right) + \cosh\left(\frac{H}{L(\omega, T)}\right) - \exp(-bi.H)} \quad (16)$$

$$Sb2(\omega, T) = -\frac{D(\omega, T)}{L(\omega, T)} \cdot \tanh\left(\frac{H}{L(\omega, T)}\right) \quad (17)$$

Previous works have investigated at the second solution (Eq. 17) [27], [28]. Our study will deal with this second solution whose module and phase are represented versus logarithm of the modulation frequency by figures 6 and 7, for different temperature values.

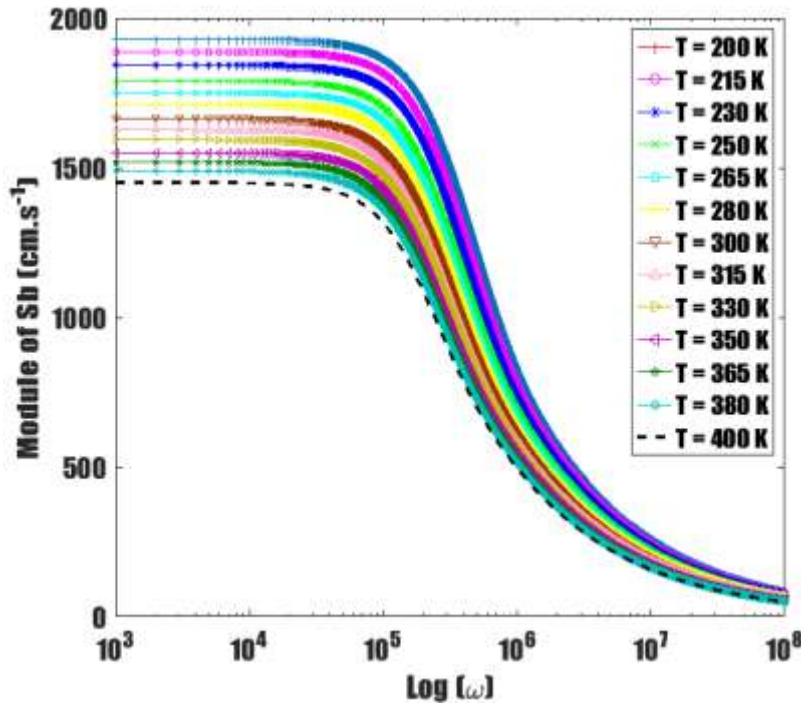


Figure 10:- Module of Sb versus frequency for different temperatures.

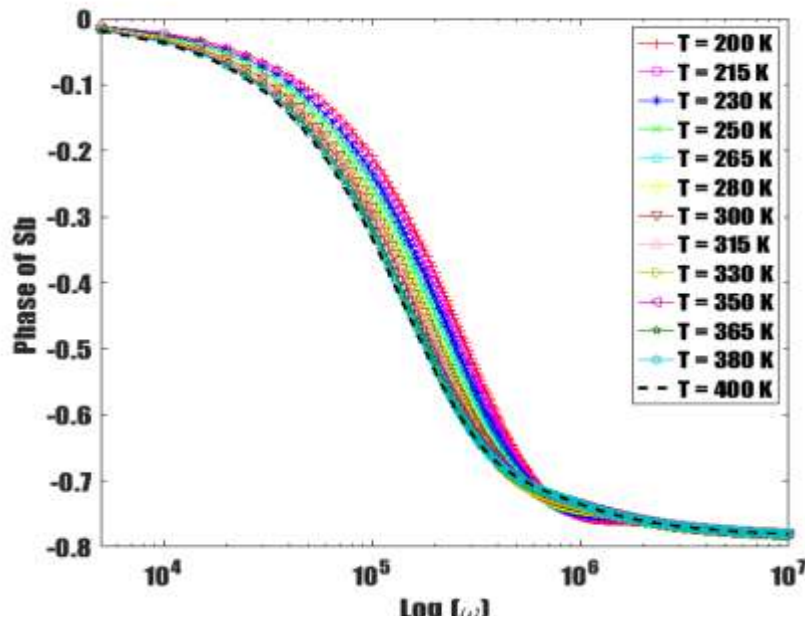


Figure 11:- Phase of Sb versus frequency for different temperatures.

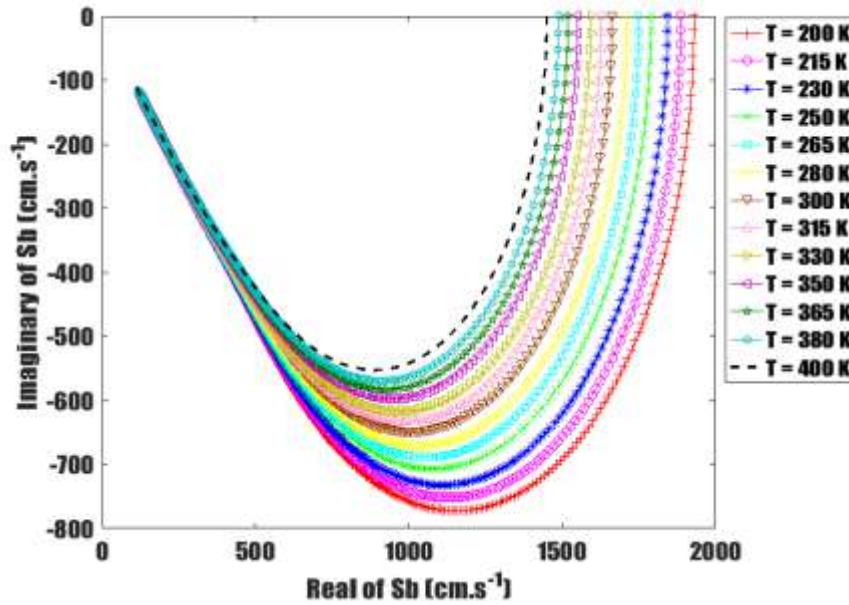


Figure 12:- Imaginary component versus real component of Sb for different temperatures.

Ac Sb in complex form (real and imaginary components, with J complex number) is presented by analogy with Maxwell-Wagner-Sillars (MWS) model [43], [44], [45], [46] and can be written as:

$$Sb(\omega, T) = Sb'(\omega, T) + J \cdot Sb''(\omega, T) \tag{18}$$

Ac phase $\phi(T)$ for a given temperature, is written as:

$$\tan(\phi(\omega, T)) = \frac{Sb''(\omega, T)}{Sb'(\omega, T)} \tag{19}$$

$Sb_{\text{ampl}}(\omega, T)$ and $\phi(\omega, T)$ correspond for a given temperature T, to the amplitude and phase component of Sb. At low frequencies ($\leq 10^4$ rad/s), the stationary regime is observed and gives constant amplitudes that decrease with the temperature T (figure. 10).

Beyond the frequency ($\gg 10^4$ rad/s), the cut-off frequency (ω_c , Sb (T)) is determined for each temperature. It is noted that the cut-off frequency decreases with the temperature T, as does the amplitude (Sb_{ampl}) (See **Table. 1**).

The phase is negative for all values of the frequency and decreases slightly with the temperature T (**figure. 11**). Figure. 12 represents the Niquyst diagram of the excess minority carriers recombination velocity for different temperatures. The radius of the semi-circles decrease with the temperature (See table 2 and table 3).

Figure 10 and figure 11 show that the sb recombination velocity decreases with temperature. Indeed, when the temperature is above the optimum temperature ($T_{\text{opt}}=300$ K) [37], thermal agitation yealds to the exponential evolution of umklapp processes that predict a temperature dependence of thermal conductivity in $1/T$ [23] [37]. This decreases the mobility of excess minority carriers and results in a decrease in the effective diffusion coefficient [23], [35] which increases recombination.

The negative phase of the ac Sb recombination velocity and the determination of electrical parameters, with the Bode and Nyquist diagrams, characterizing Sb, allow to determine the equivalent electric model [1], [11], [27], [28], [44].

Table 1:- Ac Sb cut-off frequency for different temperature values.

T (K)	200	215	230	250	265	280	300	315	330	350	365	380	400
ω_c (10^5 rad/s)	1.58	1.52	1.35	1.31	1.21	1.16	1.02	1.01	$9.5 \cdot 10^4$	$9.4 \cdot 10^4$	$9 \cdot 10^4$	$8.7 \cdot 10^4$	$8.5 \cdot 10^4$
Sb (T)	1784	1736	1710	1648	1619	1582	1552	1514	1487	1440	1414	1387	1349
Sb _{ampl}	1930	1887	1844	1790	1751	1713	1663	1628	1594	1551	1519	1489	1450
$\omega(10^5$ rad/s)	$1.3 \cdot 10^4$	10^4	$1.2 \cdot 10^4$	$9 \cdot 10^3$	$5 \cdot 10^3$	$3 \cdot 10^3$	$9 \cdot 10^3$	$8 \cdot 10^3$	$7 \cdot 10^3$	$4 \cdot 10^3$	$8 \cdot 10^3$	$6 \cdot 10^3$	$8 \cdot 10^3$

Table 2:- Maximum amplitude of the imaginary part of Sb for different temperature values.

T (K)	200	215	230	250	265	280	300	315	330	350	365	380	400
ImSb _{max} (cm/s ¹)	464.58	451.70	431.88	417.17	404.36	392.10	372.74	366.67	357.03	344.79	336.19	328.33	313.52

Table 3:- Maximum amplitude of real part of Sb for different temperature values.

T (K)	200	215	230	250	265	280	300	315	330	350	365	380	400
Re(Sb)(cm.s ¹)	1873.25	1832.14	1792.71	1740.71	1703.67	1667.52	1620.69	1586.17	1553.5	1512.19	1481.33	1452.35	1415.7
$\left(\frac{1}{\text{Re}(Sb(\text{cm.s}^{-1}))}\right) \times 10^{-5}$	53.38	54.58	55.78	57.45	58.70	59.97	61.70	63.04	64.37	66.13	67.51	68.85	70.64

Conclusion:-

The solar cell's ac back surface (p/p^+) recombination velocity that controls the recombination of the excess minority carriers has been determined. Thus, the spectroscopy method allowed the study of the Bode and Nyquist diagrams and extracted certain electrical parameters characterizing the equivalent electric model. The effect of temperature on back surface recombination velocity was explained by umklapp processes.

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