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Modified Atomic Orbital Theory Applied to the Study of $(3d^9 4s^3 D_{3,2,1})np$ 1 and the $(3d^9 4s^1 D_2)np$ Rydberg series of Cu-like Zn⁺. 2 3 Abstract 4 Precise values of high-lying resonance energies have been determined for the Rydberg series 5 $(3d^9 4s^3 D_{3,2,1})np$ and $(3d^9 4s^1 D_2)np$ (n ranging from 5 to 45), originating from the ground state 6 $(3d^{10} 4s^1 S_{1/2})$ of the Cu-like Zn⁺ ion. These calculations were carried out within the framework 7 of the Modified Atomic Orbital Theory (MAOT) method. The results show excellent 8 agreement with high-resolution measurements performed at ALS, as well as with calculations 9 obtained using the DARC codes 3 (Dirac Atomic R-matrix Codes) (Hinojosa et al., MNRAS, 10 470, 4048 (2017)). The analysis is based both on standard quantum defect theory and on the 11 MAOT procedure, which relies on the calculation of the effective charge. The study highlights 12 the relevance of the MAOT method in complementing experimental data, particularly for the 13 accurate identification of narrow resonance energies affected by overlapping spectral 14 peaks. Finally, new accurate resonance energies for high-lying levels ($n = 16 - 45$) are 15 proposed as benchmark data to 6 aid in the interpretation of spectra from trans-iron elements in 16 astrophysical environments.

17 18 Keywords 19 Photoionization, Resonance Energies, Quantum Defect, Modified Atomic Orbital Theory, 20 Rydberg Series, Synchrotron Radiation 21 22 1. Introduction 23 Resonant photoionization (PI) is a fundamental process in which a photon excites an atom or 24 ion to an unstable intermediate electronic state (autoionizing), which subsequently decays by 25 emitting a free electron. This mechanism leads to a strong enhancement of photoionization 26 cross sections and gives rise to characteristic resonant structures in the spectra. 27 In astrophysical plasmas, such as stellar atmospheres and interstellar media enriched in heavy 28 metals, zinc PI plays a key role in the ionization balance and in determining chemical 29 abundances. Zinc, particularly in the form of Zn II and Zn III ions, is used as an astrophysical 30 tracer since it is only weakly affected by depletion onto interstellar dust. Thus, the study of 31 resonant photoionization of zinc provides constraints on the chemical composition and 32 evolution of stars and galaxies based on UV observations (e.g., with HST and FUSE) (Ferland 33 et al., 1998; Kallman & Palmeri,

In laboratory plasmas, resonant photoionization of zinc is investigated using synchrotron light 35 sources (ALS, SOLEIL, PETRA III), allowing high-resolution measurements of PI cross 36 sections for different ions (Zn II, Zn III, Zn IV). These experimental measurements provide 37 essential benchmark data to test and validate advanced theoretical models (R-matrix, Dirac 38 Atomic R-matrix Codes – DARC, and semi-empirical methods such as SCUNC) (Peart et al., 39 1987; Hinojosa et al., 2017; Sakho, 2017, 2018). These studies are also applied to plasma 40 diagnostics and the production of highly charged ions in laboratory physics. 41 Although high-resolution ALS measurements have been conducted and compared with 42 advanced DARC calculations, many spectral peaks of Zn^+ ions remain overlapping, and the 43 corresponding resonances have not been clearly identified.

1 **The Modified Atomic Orbital** 44 Theory(MAOT) method (Diallo et al. 2018,2025; Sakho, 2011) has proven to be a highly 45 suitable semi-empirical method for accurately reproducing experimental photoionization data. 46 The aim of the present study is to clarify the overlapping ALS lines by using the MAOT 47 method to calculate resonance energies of the $(3d^9 4s D_3)np$, $(3d^9 4s {}^3D_2)np$, $(3d^9 4s {}^3D_1)np$, 48 and $(3d^9 4s {}^1D_2)np$ Rydberg series of Zn^+ ions. Moreover, high-lying Rydberg series provide 49 valuable data for the NIST database, where many resonances are tabulated up to $n = 60$ for 50 atomic systems such as Mg. This work also aims to extend the ALS measurements (Hinojoh 51 et al., 2017) and associated DARC calculations to higher principal quantum numbers, $n = 52$ 16 – 45. 53 The analysis of the present results is carried out within the framework of standard quantum54 defect theory and the MAOT procedure, through calculation of the effective charge. The 55 organization of this paper is as follows. Section 2 provides a brief overview of the MAOT 56 formalism, while Section 3 presents the results. The study is concluded in Section 4. 57 58 2. Theory 59 2.1. Brief description of the MAOT formalism 60 4 **In the framework of the modified atomic orbital theory (MAOT)**, the total energy of a (vl)61 given orbital is expressed in the form in Rydberg units (Diallo et al., 2018,2025). 62 □ □ □ □ 2 2

$$E_n = E_\infty - \frac{R}{(n - \delta)^2} \quad (1)$$

The details of this Equation 1 are given in our previous papers (Diallo et al., 2025). In the photoionization study, energy resonances E_n are generally measured relatively to the E_∞

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converging limit of a given $(2S+1LJ)nl$ - Rydberg series. General expression of the resonance energy E_n is given by (in Rydberg units):

$$E_n = E_\infty - \frac{R}{(n - \delta)^2} \quad (2)$$

In this equation (2), m and q ($m < q$) denote the principal quantum numbers of the $(2S+1LJ)nl$ Rydberg series of the considered atomic system used in the empirical determination of the screening constants, s represents the spin of the nl -electron ($s = 1/2$), E_∞ is the energy value of the series limit generally determined from NIST atomic database, Z represents the nuclear charge of the considered element and k is a corrective term introduced to stabilize expression (2). The only problem that one may face by using the MAOT formalism is linked to the determination of the δ -term.

The correct expression of this term is determined iteratively by imposing the general Eq. (2) to provide accurate data with nearly constant or slightly varying quantum defect values along all considered series as δ increases.

The value of δ is fixed to 1 and or 2 during the iteration.

2.2. Quantum Defect of MOAT

In general, resonance energies are analyzed from the standard quantum-defect expansion formula

$$E_n = E_\infty - \frac{R}{(n - \delta)^2} \quad (3)$$

In this equation, $R=13,60569\text{eV}$ is the Rydberg constant, E_∞ denotes the converging limit, Z_{core} represents the electric charge of the core ion, and δ means the quantum defect. From (3), the quantum defect δ is calculated as follows:

$$\delta = n - \sqrt{\frac{R}{E_\infty - E_n}} \quad (4)$$

In addition, introducing the effective nuclear charge Z^* , Eq. (2) can be rewritten in the form of Eq. (3) as follows

$$R_n Z E_{n,2} = E_{\infty} - \frac{R_{\infty}}{n^2} (Z - \sigma_1)^2 \quad (5)$$

With Z^* given by $Z^* = Z - \sigma_1 - \sigma_2$, the resonance energies are given by the formula

$$E_n = E_{\infty} - \frac{R_{\infty}}{n^2} (Z - \sigma_1 - \sigma_2)^2 \quad (6)$$

When n tends toward infinity, Eq. (6) tends to the limit $E_n \rightarrow E_{\infty} - \frac{R_{\infty}}{n^2} (Z - \sigma_1)^2$

where Z_{core} is the effective core charge. In the present work, for all the Rydberg series investigated for both Zn II, the resonance energies are given by the formula

$$E_n = E_{\infty} - \frac{R_{\infty}}{n^2} (Z - \sigma_1 - \sigma_2)^2 \quad (8)$$

The σ_2 screening constants in Equations (6) and (8) were determined empirically, based on the experimental data reported by Hinojosa et al. (2017). The corresponding values are presented in the table caption.

In contrast, the σ_2 screening constant was obtained theoretically from the simple relation $\sigma_1 = Z - Z_{core}$, where the effective core charge (Z_{core}) is directly deduced from the single photoionization process of a given XP^+ plasma ion.

$h\nu + XP^+ \rightarrow X^{P+1} + e^- \Rightarrow Z_{core} = P + 1$ So, for Cu-like Zn II ion, we find $h\nu + Zn^+ \rightarrow Zn^{2+} + e^-; Z_{core} = 2 \Rightarrow \sigma_1 = 28$

The resonance energies of the I ($3d^9 4s^3 D_3$) np , II ($3d^9 4s^3 D_2$) np , III ($3d^9 4s^3 D_1$) np , IV ($3d^9 4s^1 D_2$) np , and V ($3d^9 4s^3 D_3$) np Rydberg series,

originating from the ($3d^9 4s$) $^1S_{1/2}$ ground state of Zn^+ ions, are reported in Tables 1–5. The present MAOT calculations are compared with results obtained using the SCUNC method (Badiane et al., 2019) as well as with ALS experimental data (Hinojosa et al., 2017). In the ALS studies, resonance energies were determined only up to $n = 15$ (see Table 4). Generally, because of configuration interaction and electron–electron correlation effects, cross-section peaks tend to overlap, making the identification of spectral lines increasingly difficult as n increases. Nevertheless, the MAOT approach proves sufficiently stable to tabulate very high-lying resonances up to

identification of spectral lines increasingly difficult as n increases. Nevertheless, the MAOT approach proves sufficiently stable to tabulate very high-lying resonances up to

$n = 45$, with 118 the quantum defect remaining nearly constant across all the series investigated. Furthermore, 119 for several resonances, uncertain experimental entries indicated in parentheses are clarified. 120 Table 1 presents the quantum defects and resonance energies of the $(3d^9 4s \ ^3D_3)_{np}$ Rydberg 121 **5 series, originating from the** $3d^{10} 4s \ ^1S_{1/2}$ ground state of Zn^+ and converging to the $(3d^9 4s \ ^3D_3)$ 122 series limit of Zn^{2+} . Four uncertain ALS values are reported for the 7p, 9p, 11p, and 12p levels 123 at [26.093 eV], [26.773 eV], [27.087 eV], and [27.184 eV], respectively. These are compared 124 with the present MAOT predictions of 26.097 eV, 26.778 eV, 27.092 eV, and 27.189 eV, and 125 with the SCUNC results of Badiane et al. (2019), which give 26.102 eV, 26.782 eV, 27.095 126 eV, and 27.192 eV. The maximum discrepancy between ALS data and the new MAOT 127 calculations is 0.005 eV, while it rises to 0.009 eV in the values of Badiane et al. (2019). For 128 $n = 6$ and 8, the MAOT energies (25.409 eV and 26.510 eV) show excellent agreement with 129 ALS measurements (25.408 eV and 26.512 eV; Hinojoha et al., 2017). Across all resonances, 130 the MAOT quantum defect remains nearly constant, ranging from 1.1 to 1.2, consistent with 131 the ALS average value of 1.11 ± 0.2 (Hinojoha et al., 2017). Tables 2 and 3 lists resonance 132 **5 energies and quantum defects of the** Rydberg series $(3d^9 4s \ ^3D_2)_{np}$ and $(3d^9 4s \ ^3D_1)_{np}$, 133 originating from the ground state $(3d^{10} 4s) \ ^1S_{1/2}$ of Zn^+ ions and converging respectively to the 134 $(3d^9 4s \ ^3D_2)$ and $(3d^9 4s \ ^3D_1)$ series limits of Zn^{2+} . The results obtained with the MAOT method 135 are compared with ALS experimental measurements (Hinojoha et al., 2017) as well as with 136 the SCUNC predictions (Badiane et al., 2019). For the $(3d^9 4s \ ^3D_2)_{np}$ and $(3d^9 4s \ ^3D_1)_{np}$ series, 137 uncertain ALS values were identified at the 10p and 9p states, with corresponding energies of 138 [27.104 eV] and [27.116 eV]. The MAOT predictions for these states are 27.108 eV and 139 27.121 eV, yielding deviations of only 0.004 eV and 0.005 eV, respectively, from the ALS 140 data. For the remaining resonances between $n = 5$ and $n = 10$, the agreement between theory 141 and experiment is excellent, with energy differences never exceeding 0.008 eV. Moreover, the 142 quantum defect predicted by MAOT remains nearly constant, ranging between 1.1 and 1.2. 143 Finally, for the two

$(3d^9 4s \ ^3D)_{np}$ Rydberg series originating from the ground state $(3d^{10} 4s \ 144 \ ^1S_{1/2})$ of Zn^+ , the MAOT results are extended up to $n = 45$, confirming an almost invariant 145 quantum defect throughout the series. In the Table 4 we have the resonance energies and 146 quantum defect of the $(3d^9 4s \ ^1D_2)_{np}$ Rydberg series, originating from the $(3d^{10} 4s \ ^1S_{1/2})$ ground 147 state of the Zn^+ ion and converging to the $(3d^9 4s \ ^1D_2)$ series limit of Zn^{2+} . The results obtained 148 using the MAOT method are compared with ALS experimental data (Hinojoha et al., 2017) as 149 well as with the theoretical predictions of the SCUNC model (Badiane et al., 2019). Three 150

characteristic resonances $(3d^9 4s \ ^1D_2)6p$, $(3d^9 4s \ ^1D_2)7p$, and $(3d^9 4s \ ^1D_2)8p$ are identified as 151 overlapping states. Their resonance energies, 26.080 eV, 26.769 eV, and 27.183 eV, 152 respectively, are in close agreement with both the MAOT predictions (26.081 eV, 26.770 eV, 153 and 27.184 eV) and those from the SCUNC model. For the $(3d^9 4s \ ^1D_2)8p$ state, the MAOT 154 value shows particularly good agreement with experiment, with a deviation of only 0.001 eV, 155 compared to 0.006 eV for the SCUNC prediction, confirming the reliability of this 156 result. Except for the $(3d^9 4s \ ^1D_2)6p$ state, which shows a shift of 0.011 eV, the maximum 157 discrepancy between the new calculations and the experimental data remains below 0.07 eV. 158 For levels $n = 9$ to 15, the agreement between MAOT values and ALS measurements is 159 excellent, with differences never exceeding 0.006 eV. Finally, the quantum defect of this 160 Rydberg series remains nearly constant, between 1.1 and 1.2. The Table 5 shows the resonance 161 **energies and quantum defects of** the $(3d^9 4s \ ^3D_3)_{np}$ Rydberg series, originating from the 162 $(3d^{10} 4s \ ^1S_{1/2})$ ground state of the Zn^+ ion and converging to the $(3d^9 4s \ ^3D_3)$ series limit of Zn^{2+} . 163 The results obtained using the MAOT method are compared with the experimental ALS 164 measurements (Hinojoha et al., 2017). For this series, two overlapping resonances were 165 identified, corresponding to the 8p and 9p states, located at 26.340 eV and 26.714 eV, 166 respectively. The MAOT predictions for these levels are 26.404 eV and 26.711 eV, yielding a 167 maximum energy difference of 0.064 eV. It is worth noting that

the SCUNC method (Badiane et al., 2019) encountered difficulties in providing reliable data for this series, highlighting the relevance of the MAOT predictions as a reference for future high-resolution measurements. Finally, for all resonances considered, the quantum defect remains nearly constant, within the range 1.4-1.5. Tables 6 and 7 compare resonance energies of the $(3d^9 4s \ ^3D_3)np$ and $(3d^9 4s \ ^3D_2)np$ series, as well as the $(3d^9 4s \ ^3D_1)np$ and $(3d^9 4s \ ^1D_2)np$ series, all originating from the $3d^{10} 4s \ ^1S_{1/2}$ ground state of Zn^+ and converging respectively to the $(3d^9 4s \ ^3D_3)$ and $(3d^9 4s \ ^3D_2)$ limits of Zn^{2+} . The present MAOT results are compared with ALS measurements, the Dirac-Coulomb R-matrix (DARC) calculations of Hinojosa et al. (2017), and the pioneering merged-beams experiments of Peart et al. (1987) conducted at the Daresbury synchrotron facility. The comparison indicates that MAOT predictions agree more closely with ALS data than other theoretical approaches. The maximum relative deviation is only 0.04% for the $(3d^9 4s \ ^3D_3)7p$ state, mainly due to experimental uncertainties in the ALS value [26.093 eV]. In contrast, the DARC method shows a deviation of 0.07% for the same state, while the DARC approach exhibits a larger deviation of 0.24% at the $(3d^9 4s \ ^3D_2)7p$ level.

Tableau 1: Resonance energies (E) and quantum defect (δ) of the $(3d^9 4s \ ^3D_3)np$ Rydberg series I originate from the $(3d^{10} 4s) \ ^1S_{1/2}$ ground state of the Zn^+ ions converging to the $(3d^9 4s \ ^3D_3)$ series limit of Zn^{2+} . The present MAOT results are compared with the SCUNC results (Badiane et al., 2017) and the ALS experimental data (Hinojosa et al., 2017). The ALS energies are calibrated to $\pm 0,017 \text{ eV}$. $\sigma_1(2D5/2) = -0.770 \pm 0.048$; $\sigma_2(2D5/2) = 28.00$.

n	l	$(3d^9 4s \ ^3D_3)np$	MAOT	SCUNC	ALS	MAOT	SCUNC
\square	E	E	$ \Delta E $	$ \Delta E $			
5	1,1	24,141	24,141	24,141	0.000	0.000	6
1,1	25,409	25,415	25,408	0.001	0.007	7	1,1
26,097	26,102	[26,093]	[0.004]	[0.009]	8	1,1	26,510
26,515	26,512	0.002	0.003	9	1,1	26,778	26,782
[26,773]	[0.005]	[0.009]	10	1,1	26,961	26,965	26,966
0.005	0.001	11	1,1	27,092	27,096	[27,087]	0.005
[0.008]	[0.008]	12	1,1	27,189	27,192	[27,184]	0.005
[0.008]	[0.008]	13	1,1	27,262	27,265	(27,253)	(0.009)
(0.012)	(0.012)	14	1,1	27,320			

27,321 15 1,1 27,365 27,367 16 1,1 27,401 27,403 17 1,1 27,431 27,433 18 1,1
 27,456 27,457 19 1,1 27,477 27,478 20 1,1 27,494 27,495 21 1,1 27,509 27,510 22
 1,1 27,522 27,523 23 1,1 27,533 27,534 24 1,1 27,543 27,543 25 1,1 27,551
 27,552 26 1,2 27,559 27,559 27 1,2 27,566 27,566 28 1,2 27,571 27,572 29 1,2
 27,577 27,577 30 1,2 27,582 27,582 31 1,2 27,586 27,586 32 1,2 27,590 27,590 33
 1,2 27,593 27,594 34 1,2 27,596 27,597 35 1,2 27,599 27,600 36 1,2 27,602
 27,602 37 1,2 27,605 27,605 38 1,2 27,607 27,607 39 1,2 27,609 27,609 40 1,2
 27,611 27,611 41 1,2 27,613 - 42 1,2 27,614 - 43 1,2 27,616 - 44 1,2 27,617 -
 45 1,2 27,619 -

Tableau 2: Resonance energies (E) and quantum defect (δ) of the (3d94s 3D2)*np* Rydberg 188 series II originate from the (3d104s) 1S1/2 ground state of the Zn+ ions converging to the (3d94s 189 3D2) series limit of Zn2+. The present MAOT results are compared with the SCUNC results 190 (Badiane et al., 2017) and the ALS experimental data (Hinojoha et al., 2017). The ALS 191 energies are calibrated to $\pm 0,017$ eV. 192

<i>n</i>	II(3d94s3D2) <i>np</i>	MAOT	SCUNC	ALS	MAOT	SCUNC	δ	E	E	E	$ \Delta E $	$ \Delta E $			
5	1,1	24,294	24,294	24,294	0.000	0.000	6	1,1	25,559	25,564	25,569	0.010	0.005		
7	1,1	26,245	26,250	26,243	0.002	0.007	8	1,1	26,658	26,662	26,656	0.002	0.006		
9	1,1	26,925	26,926	0.001	0.003	10	1,1	27,108	27,112	[27,104]	[0.004]	[0.008]			
11	1,1	27,239	27,242	12	1,1	27,335	27,338	13	1,1	27,409	27,411	14	1,1	27,466	27,468
15	1,1	27,511	27,513	16	1,1	27,547	27,549	17	1,1	27,577	27,579	18	1,1	27,602	27,603
19	1,1	27,623	27,623	20	1,1	27,640	27,641	21	1,1	27,655	27,656	22	1,1	27,668	27,669
23	1,1	27,679	27,680	24	1,1	27,689	27,690	25	1,1	27,697	27,698	26	1,2	27,705	27,705
27	1,2	27,712	27,712	28	1,2	27,717	27,718	29	1,2	27,723	27,723	30	1,2	27,728	27,728
31	1,2	27,732	27,732	32	1,2	27,736	27,736	33	1,2	27,739	27,740	34	1,2	27,743	27,743
35	1,2	27,745	27,746	36	1,2	27,748	27,748	37	1,2	27,751	27,751	38	1,2	27,753	27,753
39	1,2	27,755	27,755	40	1,2	27,757	27,757	41	1,2	27,759	-	42	1,2	27,760	-
43	1,2	27,762	-	44	1,2	27,763	-	45	1,2	27,765	-				

27,451 27,450 0.000 0.001 10 1,1 27,633 27,634 27,632 0.001 0.002 11 1,1 27,764 27,765
 27,767 0.003 0.002 12 1,1 27,860 27,861 27,863 0.003 0.002 13 1,1 27,933 27,934 27,939
 0.006 0.005 14 1,1 27,990 27,991 27,994 0.004 0.003 15 1,1 28,035 28,036 28,041 0.006
 0.005 16 1,1 28,072 28,072 17 1,1 28,102 28,102 18 1,1 28,126 28,126 19 1,1
 28,147 28,147 20 1,1 28,164 28,164 21 1,1 28,179 28,179 22 1,1 28,192 28,192 23
 1,1 28,203 28,203 24 1,1 28,213 28,213 25 1,1 28,221 28,221 26 1,2 28,229
 28,229 27 1,2 28,236 28,236 28 1,2 28,242 28,242 29 1,2 28,247 28,247 30 1,2
 28,252 28,252 31 1,2 28,256 28,256 32 1,2 28,260 28,260 33 1,2 28,263 28,263 34
 1,2 28,267 28,267 35 1,2 28,269 28,269 36 1,2 28,272 28,272 37 1,2 28,275
 28,275 38 1,2 28,277 28,277 39 1,2 28,279 28,279 40 1,2 28,281 28,281 41 1,2
 28,283 - 42 1,2 28,284 - 43 1,2 28,286 - 44 1,2 28,287 - 45 1,2 28,289 -

Tableau 5: Resonance energies (E) and quantum defect (δ) of the $(3d94s\ 3D3)np$ Rydberg 203 series V originate from the $(3d104s)\ 1S_{1/2}$ ground state of the Zn^+ ions converging to the $(3d94s\ 204\ 3D3)$ series limit of Zn^{2+} . The present MAOT results are compared with the SCUNC results 205 (Badiane et al., 2017) and the ALS experimental data (Hinojoha et al., 2017). The ALS 206 energies are calibrated to $\pm 0,017$ eV.

n	MAOT	ALS	$ \Delta E $	δ	E	E											
6	1,4	25,071	25,071	0,000	7	1,4	25,917	25,910									
8	1,4	26,404	(26,340)	(0,064)	9	1,4	26,711	(26,714)	(0,003)	10	1,4	26,916	11	1,4	27,061		
12	1,4	27,166	13	1,4	27,245	14	1,4	27,306	15	1,4	27,354	16	1,4	27,393	17	1,4	27,424
18	1,4	27,450	19	1,4	27,472	20	1,4	27,490	21	1,4	27,505	22	1,4	27,519	23	1,4	27,530
24	1,5	27,540	25	1,5	27,549	26	1,5	27,557	27	1,5	27,564	28	1,5	27,570	29	1,5	27,575
30	1,5	27,580	31	1,5	27,585	32	1,5	27,589	33	1,5	27,592	34	1,5	27,596	35	1,5	27,599
36	1,5	27,601	37	1,5	27,604	38	1,5	27,606	39	1,5	27,608	40	1,5	27,610	41	1,5	27,612
42	1,5	27,614	43	1,5	27,615	44	1,5	27,617	45	1,5	27,618	208					

Tableau 6: Comparison of resonance energies of the $(3d94s\ 3D3)np$ and $(3d94s\ 3D2)np$

209 Rydberg series originating from the $(3d^{10}4s) 1S_{1/2}$ ground state of the Zn^+ ions converging to 210 the $(3d^9 4s^3 D_3)$ and $(3d^9 4s^3 D_2)$ series limits respectively of Zn^{2+} . 211 $n (3d^9 4s^3 D_3)_{np}$ Rydberg series $(3d^9 4s^3 D_2)_{np}$ Rydberg series MAOT E DARC E ALS E DAR E MAOT E DARC E ALS E DAR E 5 24,141 24,158 24,141 24,172 24,294 24,312 24,294 24,329 6 25,409 25,418 25,408 25,468 25,559 25,577 25,569 25,616 7 26,097 26,102 [26,093] 26,150 26,245 26,251 26,243 26,305 8 26,510 26,512 26,658 26,656 9 26,778 [26,773] 26,925 26,926 10 26,961 26,966 27,108 [27,104] 11 27,092 [27,087] 27,239 12 27,189 [27,184] 27,335 13 27,262 (27,253) 27,409 MAOT, Modified Atomic Orbital Theory, present calculations. 212 3 DARC, Dirac Atomic R-matrix Codes, (Hinojoha et al., 2017). 213 ALS, experimental results, (Hinojoha et al., 2017). 214 DAR, DARES BURY, experimental results (Peart et al., 1987). 215 216 217 Tableau 7: Comparison of 1 resonance energies of the $(3d^9 4s^3 D_1)_{np}$ and $(3d^9 4s^1 D_2)_{np}$ 218 Rydberg series originating from the $(3d^{10} 4s) 1S_{1/2}$ ground state of the Zn^+ ions converging to 219 the $(3d^9 4s^3 D_1)$ and $(3d^9 4s^1 D_2)$ series limits respectively of Zn^{2+} . 220 $n (3d^9 4s^3 D_1)_{np}$ Rydberg series $(3d^9 4s^1 D_2)_{np}$ Rydberg series MAOT E DARC E ALS E DAR E MAOT E DARC E ALS E DAR E 5 24,488 24,504 24,488 24,509 24,804 24,824 24,804 24,839 6 25,754 25,762 25,752 25,786 26,080 [26,069] 7 26,440 26,439 26,769 [26,773] 8 26,853 26,858 27,183 [27,184] 9 27,121 [27,116] 27,450 27,450 10 27,304 27,633 27,632 11 27,435 27,764 27,767 12 27,531 27,860 27,863 13 27,605 27,933 27,939 14 27,662 27,990 27,994 15 27,707 28,035 28,041 MAOT, Modified Atomic Orbital Theory, present calculations. 221 3 DARC, Dirac Atomic R-matrix Codes, (Hinojoha et al., 2017). 222 ALS, experimental results, (Hinojoha et al., 2017). 223 DAR, DARES BURY, experimental results (Peart et al., 1987). 224 225 226

4. Conclusion 227 1 The Modified Atomic Orbital Theory (MAOT) formalism has been applied to the 228 photoionization study of the $(3d^9 4s^3 D_{3, 2, 1})_{np}$ and $(3d^9 4s^1 D_2)_{np}$ ($n = 5$

– 45) Rydberg 229 series, originating from the $(3d^{10}4s\ ^1S_{1/2})$ ground state of the Zn^+ ion. Overall, the results 230 obtained show excellent agreement with high-resolution ALS measurements, as well as with 231 Dirac–Coulomb R-matrix calculations [(Hinojosa et al., 2017), MNRAS, 000, 1] and 232 theoretical predictions from the SCUNC method (Badiane et al., 2019). This work 233 demonstrates the ability of the MAOT approach to support experiments in resolving very 234 closely spaced resonance energies, particularly in the presence of overlapping peaks. The 235 newly determined resonance energies, obtained with high precision, are proposed as reference 236 data for the interpretation of Zn^+ spectra observed in nebulae. Furthermore, these high-lying 237 Rydberg series provide a valuable contribution to the enrichment of the NIST atomic 238 database. 239 FUNDING: 240 This research received no funding. 241 DATA AVAILABILITY 242 All data generated or analyzed during this study are included in this published article. 243 CONFLICT OF INTEREST 244 The authors of this work declare that they have no conflicts of interest 245

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