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Superposed epoch analysis of Magnetospheric Convection Electric 1 Field during geomagnetic storms due to High Speed Solar Wind 2 during solar cycle 24 3 4 5 Abstract 6 The present study concerns the time variation of the MCEF with respect to geomagnetic 7 phases as shown by SymH time variation by means of a superposed epoch analysis during 8 solar cycle 24 for HSSW. For this work the magnetosphere input energy behavior is also 9 highlighted. During IP and MP, the main effect of storm is the drop of MCEF and the 10 decreasing particles density especially during the storm MP. During MP and RP, there is 11 convection into a magnetosphere and the time for the total releasing of magnetosphere energy 12 is longer than devoted to the RP. From IP to MP + RP, are pointed out the substorm onsets and 13 the presence of substorms during the MP. From IP to MP + RP, the substorms that occur 14 during the MP and the generated storms are not HILDCAA. 15 16 Keywords 17 Superposed epoch analysis, Magnetospheric convection electric fields, solar wind high speed, 18 geomagnetic storm, geomagnetic indices 19 20 Introduction 21 22 Geomagnetic storm can be defined as a major disturbance of Earth's magnetosphere due to an 23 efficient exchange of energy from the solar wind into the space environment surrounding 24 Earth. The primary causes of geomagnetic storms at Solar are Coronal Mass Ejections 25 (CMEs) and High-Speed Solar Wind streams (HSSW). According to Gonzalez (1994) the 26 primary sources of geomagnetic storms at Earth are strong dawn-to-dusk electric fields 27 associated with the passage of southward directed Interplanetary Magnetic Fields (IMF), Bz, 28 past the Earth for sufficiently long intervals of time. 29

When CMEs and/or HSSW interact with Earth's magnetic field, they can provoke 30 significant changes in the magnetosphere with various effects on Earth (McPherron et al., 31 2008). During this disturbance, MCEF plays a key role. Many studies (e.g. Akhavan-Tafti et al., 2023; Alqeeq et al., 2025; Kim et al., 2025; Bazie et al., 2025 and Desta et al., 2026) are 33 performed to study magnetosphere variability during geomagnetic storms. Storms caused by 34 CMEs have been particularly well-studied (e.g. Gopalswamy

N., 2002; Pulkkinen et al., 2007; 35 Gopalswamy N., 2009; Benacquista et al., 2010; Ontiveros et al., 2010; Liou et al., 2016; 36 Kabore et al., 2018; Pedersen et al., 2022 and Bazie et al. 2025) but those due to high-speed 37 solar winds are less. Several scientists (e.g. Denton et al., 2006, Borovsky et al., 38 2006), Hutchinson et al., 2011) and Grandin et al., 2019) showed that HSSW storms are 39 different from CMEs storms. It is therefore important to carry out analyses specific to these 40 storms. The present study focuses on 20 storms caused by high-speed winds and/or CIRs 41 during the declining phase of solar cycle 24. In this study, we attempt to answer the following 42 questions: How does the convective electric field vary during the different phases of storms 43 caused by high-speed solar winds? What are the solar wind parameters that determine the 44 state of the magnetosphere during each phase of a storm caused by high-speed solar winds? 45 To answer these questions, we use a statistical storm-phase approach. Several studies have 46 shown that the phase-based approach is useful for understanding the processes involved 47 during magnetic storms (Wang et al. 2024; Mishra et al. 2024; Ahmed et al. 2024). However, 48 the study of storms presents a major complexity: the events do not have the same durations, 49 and the phases of the storm do not have the same characteristics. In order to standardize them 50 and extract average trends for each phase, we use a superposed epoch analysis with time 51 normalization. It is important to note that this method has been used to study geomagnetic 52 storms by Hutchinson et al. (2011) for the intensity of the storms.; Yermolaev et al. (2010) for 53 the average behavior of the geomagnetic storms; Keese et al. (2014) for studying the global 54 ion temperature during the ICME driven and CIR driven storms; Katsavrias et al. (2019) for 55 the acceleration and loss of relativistic electron. To improve the accurate behavior of the 56 storm Yokoyama and Kamide (1997) and Manu et al. (2023) used Double superposed epoch 57 analysis. 58 For us, the superposed epoch analysis allows on one hand, to equalize storms over the same 59 normalized time intervals, and also to deduce the mean, median, as well as the lower and 60 upper quartiles of the different parameters for each normalized time. These values allow us to 61 compare the variation 1 of the

and to extract the average variation during storms in order to answer the first question. To  
63 answer the second question, we analyze the relationship between solar wind  
parameters and 64 the MCEF throughout the different phases of storms. In the following,  
we will describe the 65 data and the methods and we will present the results. 66 1. Data  
and methodology 67 1.1 Data 68 Magnetospheric activities, geomagnetic storms,  
substorms are due to the energy injected into 69 the magnetosphere from the solar wind  
(Akasofu, 1981). This energy injected is responsible 70 for the dynamics of the  
magnetosphere-ionosphere system (Newell et al., 2007). To determine 71 the energy input,  
several energy coupling functions are used (Perreault and Akasofu, 1978; 72 Kan and Lee,  
1979; Murayama, 1982; Vasyliunas et al. 1982; Wygant et al. 1983; Bargatze et 73 al.,  
1986; Xu and Shi, 1986; Mac-Mahon and Gonzalez, 1997; Stamper et al., 1999; Koskinen  
74 and Tanskanen, 2002; Vichare et al., 2005; Finch and Lockwood, 2007; Wang et al.,  
2014). In 75 this work the coupling function of Wang et al. (2014) is used. 76 The energy  
input can be predicted by means of  $B_z$  component of the IMF, the solar wind 77 velocity ( $V$ )  
and the dynamic pressure ( $p$ ) (Dungey, 1961; Kan and Lee, 1979; Wygant et al., 78 1983;  
Scurry and Russell, 1991; Temerin and Li, 2006; Newell et al., 2007). It can be also 79  
noted that the magnetospheric state variables can be determined by using  $B_z$ . 80  
Therefore, in the present study, the following data are used: solar wind velocity ( $V$ ); the  $Z$   
81 component of Interplanetary magnetic Field (IM ( $B_z$ )) and the  $Y$  component of  
Interplanetary 82 Electric Field (IEF) ( $E_y$ ), solar wind particles density ( $n$ ). 83 To determine  
storm time phases, we utilized the geomagnetic index  $SymH$  (according to 84 Iyemori  
(1990),  $SymH$  measures the activity of ring current with 1 minute as time resolution 85 and  
is same as  $Dst$  with 1 hour as time resolution. To estimate the overall intensity of auroral  
86 currents the Auroral Electrojet (AE) is used in this work. 87 This study focuses on the  
variation of the magnetospheric convection electric field (MCEF) 88 under the influence of  
High-Speed Solar Wind (HSSW). It is important to note that this 89 influence can be

interpreted in terms of, on the one hand, energy injected into the 90 magnetosphere as a function of the characteristics of solar wind parameters and, on the other 91 hand, geomagnetic storms intensity and time variation. Therefore, we will use the geomagnetic 92 activity indices SymH and AE and solar wind parameters that are available on the 93

OMNIWEB website [https://omniweb.gsfc.nasa.gov/form/omni\\_min.html](https://omniweb.gsfc.nasa.gov/form/omni_min.html) and recorded in Table 94 1. 95 96 Table 1: Geomagnetic index and solar wind parameters 97 Parameters (Unit) Resolution Description Utilisation V (Km/s) 5 min Solar wind speed For analyzing changes in solar wind speed Bz (nT) 5 min Component along the z-axis of the interplanetary magnetic field For analyzing conditions favorable to reconnection Ey (mV/m) 5 min Component of the electric field along the y-axis of the interplanetary magnetic field For calculating the MCEF AE (nT) 5 min AE index for monitoring auroral activity For evaluating auroral activity SymH (nT) 5 min Index quantifying variations in the ring current For dividing storms into different phases n, particles density (particles/cm<sup>3</sup>) 5 min Densité des particules du vent solaire 98 1.2 Methodology 99 1.2.1 Study of geomagnetic storms 100 a) Determination of geomagnetic storms 101 In this study that covers the period from 2008 to 2018, we focus on geomagnetic storms. The 102 start of a storm is identified by the condition  $SymH < -60 \text{ nT}$  and the end by the 103 condition  $SymH > -20 \text{ nT}$ . There were 125 storms during the period in question, regardless 104 of their origin or source. 105 Several classifications of storms exist: 106 a) Classification according to Hutchinson et al. (2011) and Walach et al. (2019). 107 Low-intensity storms:  $-150 \text{ nT} < SymH_{min} \leq -80 \text{ nT}$  108 Moderate storms:  $-300 \text{ nT} < SymH_{min} \leq -150 \text{ nT}$  109

Intense or strong storms:  $SymH_{min} \leq -300 \text{ nT}$  where  $SymH_{min}$  is the minimum value of 110 SymH 111 b) Classification according to Li et al. (2010; 2012) 112 Moderate storms: 100

The storm time is divided into three phases with respect to SymH time variation (figure 1): (a) 143 the Initial Phase (IP) with time interval  $(t_0-t_1)$ ; the Main Phase (MP) characterized

by a time 144 interval ( $t_1-t_2$ ) and the Recovery Phase (RP) with time interval ( $t_2-t_3$ ) (table 2) 145 146 Table2: Start and End phases of the retained geomagnetic storms 147 Number of storm Start of IP Start of MP or End of IP Start of RP or End of MP End of RP SymHmin

Storm Number	Start of IP	Start of MP or End of IP	Start of RP or End of MP	End of RP	SymHmin
1	03/09/2008 14:50	03/09/2008 23:15	04/09/2008 03:05	04/09/2008 11:00	-67
2	10/10/2008 03:05	11/10/2008 07:20	11/10/2008 11:30	14/10/2008 22:55	-64
3	21/07/2009 03:20	21/07/2009 22:25	22/07/2009 05:55	25/07/2009 11:35	-93
4	01/05/2010 12:30	02/05/2010 08:15	02/05/2010 20:30	05/05/2010 05:35	-75
5	04/02/2011 07:05	04/02/2011 16:55	04/02/2011 21:20	07/02/2011 02:15	-67
6	01/03/2011 07:40	01/03/2011 09:20	01/03/2011 14:25	03/03/2011 11:50	-71
7	31/05/2013 13:55	01/06/2013 01:35	01/06/2013 07:45	03/06/2013 01:10	-134
8	07/12/2013 15:20	07/12/2013 22:40	08/12/2013 08:30	09/12/2013 05:50	-72
9	03/07/2015 23:10	04/07/2015 15:55	05/07/2015 04:55	07/07/2015 20:05	-86
10	05/03/2016 13:05	06/03/2016 15:00	06/03/2016 21:20	09/03/2016 03:05	-120
11	23/10/2016 02:40	23/10/2016 05:50	25/10/2016 22:55	28/10/2016 17:45	-80
12	26/03/2017 19:55	27/03/2017 01:00	27/03/2017 16:10	29/03/2017 20:45	-86
13	26/09/2017 19:55	27/09/2017 05:30	28/09/2017 05:50	29/09/2017 07:35	-74
14	06/11/2017 12:25	07/11/2017 04:35	08/11/2017 04:05	10/11/2017 01:00	-89
15	20/11/2017 07:40	20/11/2017 17:20	21/11/2017 06:50	24/11/2017 12:25	-60
16	05/05/2018 04:05	05/05/2018 14:20	06/05/2018 02:30	08/05/2018 01:35	-66

Figure 1: SymH time variation with storm time phases 150 151 1.2.2 Superposed Epoch Analysis method 152

Studying storms presents difficulties due to their different characteristics. Indeed, they have 153 different durations and intensities. This makes it difficult to calculate averages per phase for a 154 set of storms and to perform a general analysis of trends for each phase. An alternative is to 155 use the method called “superposed epoch analysis with time normalization”. For more details 156 about this method see Walton and Murphy (2022). Their method execution code can be 157 found at the following link [https://github.com/samwalton7645/SEA\\_Code](https://github.com/samwalton7645/SEA_Code). 158 To analyze storm parameters (Bz, AE,

Em, Wrecon, Wothers, n, SymH) , each storm SymH is 159 divided into three phases (see figure 1). For analyzing each storm phase effect on the 160 parameters, we define four cases that can be named events. To apply the Superposed Epoch 161 Analysis (SEA) method, firstly, the event time must be divided into two parts named in the 162 method phases. The first one, from the beginning of the event (*tbegin* ) to the event 163 development called by Walton and Murphy (2022) epoch (*tepoc* □ ) and the second one from 164 the epoch (*tepoc* □ ) to the end of the event (*tend* ); secondly, each part time interval is 165 normalized in order to transform it to a standard interval  $[0, 1]$  ; thirdly, the standard interval is 166 binned into a set of equally spaced bins; fourthly, for each part interval, a set of statistics 167 (mean, median, etc.) is then determined for the data residing in each bin. 168 a) Time normalization 169 The SEA code is run four times with respect to the two parts or code phases per event. We 170 have: 171 □ Run 1: Obtention of parameters time variation for IP and MP 172 *Code\_PPhase 1 = IP* and *code\_PPhase 2 = MP* 173 This is done for having on one hand the effects of IP and the effects of MP, on the other hand. 174 The time intervals are ( $t_0-t_1$ ) and ( $t_1-t_2$ ) 175 □ Run 2: Obtention of parameters time variation for MP and RP 176 *Code\_PPhase 1 = MP* and *Code\_PPhase 2 = RP* 177 The SEA code program is run for having on one hand, the effects of MP and the effects of RP, 178 on the other hand. The time intervals are ( $t_1-t_2$ ) and ( $t_2-t_3$ ) 179 □ Run 3: Obtention of parameters time variation for IP + MP and RP 180 *Code\_PPhase 1=IP + MP* and *Code\_PPhase 2 = RP* 181

The SEA code script is executed for having on one the hand, the global effects of IP and MP 182 together and on the other hand, the effects of RP. The time intervals are ( $t_0-t_2$ ) and ( $t_2-t_3$ ) 183 □ Run 4: Obtention of storm parameters time variation for IP and MP + RP 184 *Code\_PPhase 1=IP* and *Code\_PPhase 2 = MP + RP* 185 This is done for having on one hand, the effects of IP and on the other hand, the global effects 186 of MP and RP together. The time intervals are ( $t_0-t_1$ ) and ( $t_1-t_3$ ). 187 188 To apply SEA method, we convert each storm phase intervals into a standard one  $[0, 1]$  . The 189 process

is:having  $\forall t \in [t_i, t_j]$  with  $i = 0, 1, 2$  and  $j = 1, 2, 3$  we define the normalized time as follows:  
 $t_{norm} = \frac{t - t_i}{t_j - t_i}$ . This leads to a normalized interval so that for a given time  $t_{norm}$  we  
 191 have  $t_{norm} \in [0, 1]$  192 For each phase and for each instant  $t$ , we calculate the time  
 elapsed between the start of the 193 phase and instant  $t$ . That is, the time elapsed  
 between  $t_0$  and  $t$  for any instant  $t$  of phase 1 or 194 between  $t_1$  and  $t$  for any instant  $t$  in  
 phase 2 or for any instant  $t$  between  $t_2$  and  $t$  for any 195 instant  $t$  in phase 3. Each  
 elapsed time value is converted to a value between 0 and 1 by 196 dividing it by the total  
 duration of the phase. Thus, 0 corresponds to the beginning of the 197 phase and 1 to the  
 end of the phase. This means that regardless of the duration of the storms, 198 they can  
 be aligned on an evolution scale. 199 200 b) binning into a set of equally spaced bins 201  
 Firstly, each codephase standard interval is binned. The number of bins is given in table 3.  
 202 Secondly, the binned normalized time is obtained by starting to count bins at the  
 epoch time. 203 Thus, bin numbers before epoch time are negatives and that after epoch  
 time are positive. 204 205 Table 3: Number of bins per code phase for given Run 206 207  
 Run number Number of bins for Code\_Phase1 Number of bins for Code\_Phase2 1 40 120  
 2 80 100 3 40 100

4 40 40 208 1.2.3 Determination of Magnetospheric Convection Electric Field (MCEF)  
 209 To determine the values of the magnetospheric convection electric field (MCEF), we  
 will use 210 the relationship proposed by Wu et al. (1981) and validated by Revah and  
 Bauer (1983), 211 which links the  $E_y$  component of the Interplanetary Electric Field (IEF) to  
 $B_z$  component of 212 the Interplanetary Magnetic Field (IMF):  $E_M = 0.13 E_y + 0.09$  with  $E_y$   
 $= -VB_z$  where  $V$  is 213 the solar wind speed. 214 215 1.2.4 Determination of the input  
 energy components 216 To determine the energy transferred during the coupling between  
 the solar wind and the 217 magnetosphere, we will use the relationship described by Wang  
 et al. (2014): 218  $W_{in} = 3.78 \cdot 10^7 n \cdot 0.24 V^{1.47} B T^{0.86} [\sin^2 \theta / 2 + 0.25]$  where: 219  $\theta$  is  
 the shock incidence angle of the interplanetary magnetic field. It is defined as follows:  
 220  $\theta = \tan^{-1} \frac{B_y}{B_z}$  if  $B_z > 0$   $\theta = \pi - \tan^{-1} \frac{B_y}{B_z}$  if  $B_z < 0$  221  $V$  the solar wind speed,

$B_T$  the IMF magnitude expressed as:  $B_T = B_y^2 + B_z^2$ ,  $n_{sw}$  the solar wind particles density. The equation that gives the input energy can be rewritten as:  $W_{in} = 3.78 \cdot 10^7 n_{sw}^{0.24} V^{1.47} B_T^{0.86} \sin^2 \theta + 3.78 \cdot 10^7 n_{sw}^{0.24} V^{1.47} B_T^{0.86} \cdot 0.25$ . It can be seen that we have the sum of two expressions which can be expressed as:  $W_{recon} = 3.78 \cdot 10^7 n_{sw}^{0.24} V^{1.47} B_T^{0.86} \sin^2 \theta$  and  $W_{others} = 3.78 \cdot 10^7 n_{sw}^{0.24} V^{1.47} B_T^{0.86} \cdot 0.25$ . therefore  $W_{in} = W_{recon} + W_{others}$  The first term ( $W_{recon}$ ) represents the injected electromagnetic energy, mainly due to magnetic reconnection on the day side, which varies with  $\sin^2(\theta)$ . This term reaches its

maximum when the IMF is oriented towards the south ( $\theta = 180^\circ$  and  $\sin^2 \theta = 1$ ) and cancels out when the IMF is oriented towards the north ( $\theta = 0^\circ$ ). The second term ( $W_{others}$ ) represents an energy contribution independent of the clock angle, attributed to other processes such as reconnection at high latitudes (in the lobes), viscous interactions, or mechanical energy transfer (e.g., via Kelvin–Helmholtz instabilities). This component becomes relatively more important when the IMF is oriented northward.

**2. Results and Discussions**

Figures 2-6 show the bin normalized time (BNT) for some solar wind parameters (particles density, Z component of IMF), the geomagnetic indices AE and SymH and the calculated injected energy from solar wind into the magnetosphere components (reconnection energy and other sources of energy). In figures 2-5, panel a concerns AE index, panel b is devoted to Bz, panel c shows the BNT evolution of SymH, panel d concerns particles density, panel e, presents the BNT variation of MCEF and panels f and g are devoted to  $W_{recon}$  and  $W_{others}$ , respectively.

**2.1. Case of Run 1**

In figure 2 BNT varies from -40 to 39 BNT with IP (from -40 to 0 BNT) and MP (from 0 to 39 BNT). It can be seen that during IP all parameters exhibit constant value close to zero, except particles density and the injected energy due to other sources. This observation let us say before storm (1) there is no convection electric, (2) magnetosphere is submitted to input energy independently to the coupling solar wind magnetosphere interactions; (3) the

particles density is remained at 12.5 particles /cm<sup>3</sup>. When Bz decreases and tends toward zero there is a slight increase of other sources input energy and particles density. When Bz passes from northward to southward we enter into MP and we observe a sudden increase of all parameters. During the remain southward Bz, the particles density decreases and tends toward a value before the beginning of the storm. For the other parameters, they slightly increase. At the end of MP, the reconnected input energy drops from 0 W to 3 W while that of the other sources passes from 0.25 W to 1.25 W and the MCEF value goes from 0 mV/m to 0.4 mV/m.

The main effect of storm is the drop of all parameters since Bz passes southward and the smooth increase of all parameters except the particles density which highlights a maximum when Bz passes southward and after decreases. This shows the loss of particles during storm MP.

Figure 2: Solar wind parameters and geomagnetic indices BNT variation during IP and MP. From top to bottom: a) AE index; b) Bz; c) SymH; d) particles density; e) MCEF; f) input energy due to reconnection and g) input energy due to the other sources

2.2. Case of Run 2 Figure 3 shows BNT evolution of solar wind parameters and geomagnetic indices (SymH and AE) and input energy components for MP and RP. BNT varies from -40 to 0 for MP and from 0 to 100 for RP. It can be seen, when Bz remains southward, the decreasing of SymH and particles density and the increasing of the other parameters. At the minimum value of SymH (-75 nT), the end of MP and the beginning of RP, Bz reaches 0 nT, AE its maximum value (850 nT). Particles density, MCEF, and input energy curves present a maximum (15 particles/cm<sup>3</sup>, 0.5 mV/m, 3 W and 1.3 W, respectively) before the end of MP at -5 BNT. During RP, all parameters decrease and tend smoothly toward 0.8 mV/m for MCEF, 0.5 W for the

reconnecting energy and 0.25 W for that due to other sources. 286 The MCEF not reaches 0 mV/m that exhibits the remain of convection into the 287 magnetosphere, even though there is a decreasing of the input energy, it is not totally 288 dissipated after the end of storm. This pointed out that there is remaining energy at the end of 289 the storm. Consequently, the input energy is not completely dissipated. The consequence is 290 that the time for total releasing of magnetosphere energy is longer that devoted to the RP. 291 292 293 294 295

296 Figure 3: Same as Figure 2 but during MP and RP 297 298

299 2.3. Case of Run 3 300 301 Figure 3 presents two parts of storm time variation: the whole IP and MP (from -80 BNT to 0 302 BNT) as the first part and the RP (from 0 BNT to 99 BNT) as the second part. 303 During the whole IP and MP, the parameters (AE, n, Em, Wrecon, Wothers) simultaneously 304 increase and during RP they decrease at the same time and tend toward their minimum value. 305 At -40 BNT Bz becomes northward and parameters decrease. AE, MCEF and reconnection 306 input Energy are more sensitive to this change in Bz direction. When Bz passes from 307 northward to southward, AE increases until the end of MP and decreases after. Its minimum 308 value is superior to that before storm. This proves that the magnetosphere does not return to 309 its initial state after the storm has passed. The particles density decreases when Bzturns 310 southward and stabilizes at 2 particles/cm<sup>3</sup> close very lower than its value (8 particles/cm<sup>3</sup>) 311 before the storm. Everything happens as the storm depletes the magnetosphere from a particle 312 perspective. 313 The MCEF increases when Bz passes southward and starts decreasing before the end of MP. 314 The decrease continues during the recovery phase, tending towards zero and very close to the 315 state before the storm. We can conclude that the magnetospheric convection electric field 316 created by the storm fades away at the end of the storm. 317 The variability of input energy is the same as that of MCEF but with a pronounced slope at 318 the beginning of the decreasing. The final value

(0.38 W) at the end of the storm is higher 319 than that before storm ( $\sim 0$  W) for the input energy due to reconnection while the input energy 320 coming from the other sources last value ( $\sim 0.39$  W) is the same as before the storm. We can 321 conclude that the magnetosphere gains energy after the storm in terms of energy due to 322 reconnection despite tail reconnection and substorm activity allow the system to return to a 323 more normal configuration with respect to the size of the storm (Hutchinson et al., 2011; 324 Gonzalez et al., 1994; Daglis et al., 1999; Liemohn et al., 1999 and Reeves et al., 2003) 325 326

327 Figure 4: Same as Figure 2 but during IP +MP and RP 328 329 330

331 2.4. Case of Run 4 332 333 Figure 4 shows the evolution through all storm phases in accordance with the temporal 334 evolution of SymH of the parameters AE index, IMF component Bz, the density of particles in 335 the magnetosphere n, the MCEF Em, and the components of energy injected into the 336 magnetosphere, namely the reconnection component Wrecon and the component due to other 337 sources Wothers. In this figure, the IP starts from -42 bin normalized time (BNT). and ends at 0 338 BNT. The whole main and recovery phases (MP and RP) start from 0 to 120 BNT. At 0 BNT, 339 the IMFBz turns from northward to southward and reaches its minimum at 7.5 BNT. At the 340 same time (0 BNT) all parameters increase and reach their maximum at 7.5 BNT. 341 We have 610 nT, for AE; - 10 nT for Bz; 40.6 mV/m for Em, 3.75 W for Wrecon and 1 W for 342 Wothers. The particles density n reaches its maximum at 0 BNT with  $30 \text{ cm}^{-3}$ . During the RP, 343 the MCEF decreases and fluctuates around a level slightly higher than the level before the 344 storm 345 The particles density decreases and reaches a level lower than that before the storm. The 346 injected energies (Wrecon and Wothers) and AE index decrease and stabilize at levels slightly 347 higher than those before the storm. 348 The prompt increases of AE at 0 BNT when Bz turns from northward to southward lets us 349 assert that at this time there is an increase of the overall intensity of auroral current 350

(Nakamura et al., 2015) and characterizes not only a substorm onsets (Wang et al., 2014) but also the presence of substorms because these are observed to occur during the MP of magnetic storms (Gonzalez, 1994). As the sudden increases of AE index appears during the MP instead of RP and AE values are less than 1000 nT it emerges that the present storm events are not those qualified by Tsurutani and Gonzalez (1987) and Hajra et al. (2014) High-Intensity, Long-Duration, Continuous AE Activity (HILDCAA) events. The maximum AE reached during the main phase and its subsequent exponential decrease confirms the assumption made by Kamide and Fukushima (1971) that the rate of energy injection into the annular current depends on the AE index. Figure 4 exhibits that solar wind velocity increases when the particles density decreases. This expresses the behavior of storms provoked by CIRs as asserted by Hutchinson et al. (2011).

Figure 2: Solar wind parameters and geomagnetic indices BNT variation during IP and MP +RP. From top to bottom: a) Solar wind velocity; b) AE index; c) Bz; d) SymH; e)

particles density; f) MCEF; g) input energy due to reconnection and h) input energy due to the other sources

Conclusion

The MCEF time variation was studied by means of superposed epoch analysis during geomagnetic storms caused by HSSW for solar cycle 24. The analysis of MCEF through the three phases of storm time with respect to SymH variation shows that: From IP to MP, the main effect of storm is the drop of MCEF since Bz passes southward and its smooth increasing during the storm MP. From the MP to RP, the convection remains into the magnetosphere even though there is a decreasing of the input energy. This energy does not completely dissipate. From IP + MP to RP the MCEF created by the storm fades away at the end of the storm. The magnetosphere gains energy after the storm in terms of energy due to reconnection despite tail reconnection and substorm activity allow the system to return to a

381 more normal configuration. From IP to MP + RP, we have the substorm onsets during  
the 382 storm MP and the behavior of storms due to CIRs. The concerning storms are  
different to 383 those of HILDCAA events. 384 385 Bibliography 386 Ahmed O., B.  
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