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RESEARCH ARTICLE

FINITE VOLUME METHOD DISCRETIZATION OF MODIFIED NAVIER-STOKES EQUATION

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Abstract

This study has come up with a numerical scheme that arises from finite volume discretization of Modified Navier-Stokes Equation. Modified Navier-Stokes Equation in the x-z axis was coupled with continuity equation to obtain the Pressure Equation. Using the pressure equation, pressure field can be determined in each control volume on a staggered grid.

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Introduction:-

Navier-Stokes Equation is a partial differential equation representing a force of fluid flow field. Frenchman M. Navier and Englishman G. Stokes independently obtained the equation in first half of the nineteenth century and it is named to their honour. It can describe many phenomena of scientific interest. Despite its wide application, it still has challenges in its existence of smooth solutions because of its non-linearity terms and lack of explicit equation for pressure.

This study formulated a numerical solution for a Modified Navier-Stokes Equation in the x and z directions of the Cartesian plane using Finite Volume Method (FVM).

Modified Navier-Stokes Equation

Three-dimensional Navier-Stokes Equation is given in the form:

$$\left[\frac{\partial}{\partial t} (\rho \vec{v}) + \nabla \cdot (\rho \vec{v} \vec{v}) \right] = -\nabla P + \nabla \cdot (\bar{\tau}) + \rho f \quad (1)$$

The RHS (right hand side) of Navier-Stokes Equation are the sum of forces acting on a flowing fluid element. The force due to viscosity, shear stress is defined by the gradient:

$$\vec{F}_{vs} = \nabla \cdot \tau dV = \nabla \cdot \tau dx dy dz \quad (2)$$

where τ is the stress tensor.

The viscous term in the Navier-Stokes Equation is modified into a diffusion term given by:

$$\vec{F} = \nabla \cdot (\Gamma \nabla \tau_{\omega})$$

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where Γ is a diffusion coefficient and wall shear stress (τ_w)

If equation(3) replaces the viscous term in Navier-Stokes Equation, it changes to Modified Navier-Stokes Equation. Modified Navier-Stokes Equation models fluid flow on the boundary which fluid flows. This concept of boundary layer was first introduced by a German scientist Ludwig Prandtl in the year 1904. Modified Navier-Stokes Equation in x and z directions coupled with Continuity Equation in a Cartesian plane are given as follows:

x- direction

$$\frac{\partial \rho u}{\partial t} + \nabla \cdot (\rho u \vec{v}) = -\frac{\partial P}{\partial x} + \nabla \cdot (\Gamma \nabla \tau_w) \quad (4)$$

z- direction

$$\frac{\partial \rho w}{\partial t} + \nabla \cdot (\rho w \vec{v}) = -\frac{\partial P}{\partial z} + \nabla \cdot (\Gamma \nabla \tau_w) + \rho g_z \quad (5)$$

Continuity Equation

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \quad (6)$$

Continuity Equation expresses incompressibility of the fluid that is flowing.

u - is velocity component of the fluid in the x-direction, w - is velocity component of the fluid in the z-direction, ρ - density of the fluid flowing, τ_w - wall shear stress, Γ - wall shear stress diffusion coefficient.

3.0 Discretization of Modified Navier-Stokes Equation Using Finite Volume Method on a Staggered Grid

Finite Volume Method exhibits the following features:

- (i) Formulations can be based surface integral of normals that guarantee conservation properties throughout the domain.
- (ii) Complex geometries and unstructured meshes are easily accommodated, no co-ordinate transformation is needed.
- (iii) Neumann boundary conditions are easily enforced.

Finite volume method is preferred over other numerical methods because it ensures that discretization is conservative locally and globally, i.e. mass, momentum and energy are conserved in discrete sense.

In staggered grid discretization the scalars are stored at the centre of the control volume and velocity are centered at the faces of control volumes. If a uniform grid is used, the velocity locations are exactly at the midway between the grid points. Staggered grid for dependent variables was first used by Harlow and Welch in their well-known Mac (Marker and Cell) method. Since then it has been used by many researchers specifically SIMPLE (Semi-Implicit Method for Pressure Linked Equations) Procedure by Patankar and Spalding (1972).

Integration of equations(4),(5) with respect to time dependent control volume $v(t)$ is:

$$\int_{v(t)} \left[\frac{\partial \rho u}{\partial t} + \nabla \cdot (\rho u \vec{v}) = -\frac{\partial P}{\partial x} + \nabla \cdot (\Gamma \nabla \tau_w) \right] dv(t) \quad (7)$$

$$\int_{v(t)} \left[\frac{\partial \rho w}{\partial t} + \nabla \cdot (\rho w \vec{v}) = -\frac{\partial P}{\partial z} + \nabla \cdot (\Gamma \nabla \tau_w) + \rho g_z \right] dV(t) \quad (8)$$

Using divergence theorem, the volume integrals are converted into surface integrals.

Finite Volume Method discretized Modified Navier-Stokes Equation in the x-direction is given by:

$$\left[(\rho u)_P^{n+1} - (\rho u)_P^n \right] \Delta x \Delta z + (\rho u)_P^{n+1} (\Delta z)_e - (\rho u)_W^{n+1} (\Delta z)_w + (\rho w)_P^{n+1} (\Delta x)_n - (\rho w)_S^{n+1} (\Delta x)_s = D_P (\tau_w)_P^{n+1} +$$

$$\sum_{nb} D_{nb} (\tau_w)_{nb} + (P_P - P_E)^{n+1} \Delta z \quad (9)$$

$$(\rho u)^n \Delta x \Delta z - (\rho u)_P^{n+1} (\Delta z)_e + (\rho u)_W^{n+1} (\Delta z)_w - (\rho w)_P^{n+1} (\Delta x)_n + (\rho w)_S^{n+1} (\Delta x)_s + D_P (\tau_\omega)_P^{n+1} = b \tag{10}$$

where b is the mass source term. Hence equation (9) can be written as:

$$(\rho u)_P^{n+1} \Delta x \Delta z = \sum_{nb} D_{nb} (\tau_{nb}) + b + (P_P - P_E)^{n+1} \Delta z$$

$$a u_e = \sum_{nb} D_{nb}^u (\tau_\omega)_{nb} + b^u + (P_P - P_E)^{n+1} \Delta z$$

$$u_e = \frac{\sum_{nb} D_{nb}^u (\tau_\omega)_{nb} + b^u}{a} + \frac{\Delta z}{a} (P_P - P_E)$$

$$u_e = \frac{\sum_{nb} D_{nb}^u (\tau_\omega)_{nb} + b^u}{a} + d_e (P_P - P_E)^{n+1}$$

$$u_e = \tilde{u} + d_e (P_P - P_E)^{n+1} \text{ where } d_e = \frac{\Delta z}{a}, \tilde{u} = \frac{\sum_{nb} D_{nb}^u (\tau_\omega)_{nb} + b^u}{a} \tag{11}$$

u_e gives the velocity in the x-direction at (n+1) control volume face value.

Finite Volume Method discretized Modified Navier-Stokes Equation in the z-direction is given by:

$$[(\rho w)^{n+1} - (\rho w)^n] \Delta x \Delta z + (\rho w)_P^{n+1} (\Delta x)_n - (\rho w)_S^{n+1} (\Delta x)_s + (\rho w)_P^{n+1} (\Delta z)_e - (\rho w)_W^{n+1} (\Delta z)_w = (P_P - P_N)^{n+1} (\Delta x)$$

$$+ D_P (\tau_\omega)_P^{n+1} + \sum_{nb} D_{nb} (\tau_\omega)_{nb} + (\rho g_z)^{n+1} \Delta x \Delta z$$

where mass term can be written as:

$$(\rho w)^n \Delta x \Delta z + (\rho w)_S^{n+1} (\Delta x)_s - (\rho w)_P^{n+1} (\Delta x)_n + (\rho w)_W^{n+1} (\Delta z)_w - (\rho w)_P^{n+1} (\Delta z)_e + D_P (\tau_\omega)_P^{n+1} + (\rho g_z)^{n+1} \Delta x \Delta z = b \tag{12}$$

Hence the equation (12) can be rewritten as:

$$w_n = \frac{\sum_{nb} D_{nb} (\tau_\omega)_{nb} + b}{a} + \frac{\Delta x}{a} (P_P - P_N)^{n+1} \tag{13}$$

$$w_n = \tilde{w}_n + d_n (P_P - P_N)^{n+1}$$

$$\text{where: } d_n = \frac{\Delta x}{a}, \tilde{w}_n = \frac{\sum_{nb} D_{nb} (\tau_\omega)_{nb} + b^w}{a}$$

w_n is the velocity in the z-direction at (n+1) control volume face value

Discretization of Continuity Equation

In two dimension x-z axis is given by differential equation:

$$\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0 \tag{14}$$

Integration of equation (14) over control volume (CV) gives:

$$\int_{cv} \left[\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} \right] dv$$

$$(u_{i+1} - u_i)\Delta z = -(w_{j+1} - w_j)\Delta x \tag{15}$$

which is Finite Volume Method discretized form of Continuity Equation where u_{i+1}, u_i are velocities in the x-direction and w_{j+1}, w_j are velocities in the z-direction.

Pressure Equation

Substituting equation (11) and equation (13) into equation (15) gives:

$$\begin{aligned} [\tilde{u}_n + d_e(P_p - P_E) - u_i]\Delta z &= -[\tilde{w}_n + d_n(P_p - P_N) - w_i]\Delta x \\ \tilde{u}_n\Delta z + d_e\Delta z P_p - d_e\Delta z P_E - u_i\Delta z &= -\tilde{w}_n\Delta x - d_n\Delta x P_p + d_n\Delta x P_N + w_i\Delta x \\ (d_e\Delta z + d_n\Delta x)P_p - d_n\Delta x P_N - d_e\Delta z P_E + \tilde{u}_i\Delta z + \tilde{w}_i\Delta x - w_i\Delta x - u_i\Delta z &= 0 \\ a_p P_p - a_n P_N - a_e P_E + (\tilde{u}_i - u_i)\Delta z + (\tilde{w}_i - w_i)\Delta x &= 0 \end{aligned} \tag{16}$$

where $a_n = d_n\Delta x, a_e = d_e\Delta z, a_p = d_e\Delta z + d_n\Delta x$

Substitution of $\tilde{u}_i = \sum_{nb} D_{nb}^u(\tau_\omega)_{nb}^u + b^u$ and $\tilde{w}_i = \sum_{nb} D_{nb}^w(\tau_\omega)_{nb}^w + b^w$ into equation (15) and after algebraic simplification gives:

$$a_p P_p = a_e P_E + a_n P_N + b^p \tag{17}$$

where $b^p = -\Delta z \sum_{nb} D_{nb}^u(\tau_\omega)_{nb}^u - \Delta z b^u + u_i\Delta z + w_i\Delta x - \Delta x \sum_{nb} D_{nb}^w(\tau_\omega)_{nb}^w - \Delta x b^w$ which is the Pressure Equation.

Conclusion:-

The Pressure Equation is used to estimate the values pressure P(i) on a staggered grid in the x-z directions on each control volume.

flux out
w

flux influx out

x-axis

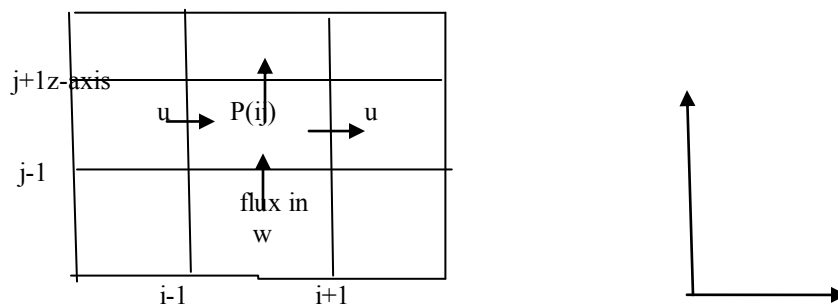


Fig 1: 3x3 mesh for a staggered grid

Thereafter horizontal velocity (u) and vertical velocity (w) at each control volume interfaces can be approximated using SIMPLER (Semi Implicit Method for Pressure Linked Equations Revised) algorithm.

Bibliography:-

1. Anderson J.D, J.F Wendt (2009) Computational Fluid Dynamics, 3rded Springer-Verlag Berlin Heidelberg, Pp20-80

2. BengtAndersson, Ronnie A., Love H., Mikael M., Rahman S.,Berend V.W (2012)Computational Fluid Dynamics for Engineers, Cambridge University Press Pp52-60.
3. Moukalled F., Mangani L., Darwish M. (2015) The Finite Volume Method inComputational Fluid Dynamics Springer Pp. 9-41,365-394,561-653,752-760
4. PatankerSuhas V. (1980) NUMERICAL HEAT TRANSFER AND FLUID FLOWTaylors & Francis Publishers Pp113-135.
- 5.Versteeg HK, Malalasekera W (1995) An Introduction to Computational Fluid Dynamics, The Finite Volume Method, Longman Group Limited Pp. 11-20
- 6.Vrushali A. Bokil, Nathan L. Gibson (2007) Finite Difference, Finite Element and Finite Volume Methods for Numerical Solutions of PDEs,DOEMultiscaleSummer School Pp. (4-40).